

Fig. 2 Effect of pressure gradient on wake thickness.

pressure. The curves at the higher altitude have the characteristic of a partially frozen flow, whereas the lower altitude curves are characteristic of chemical equilibrium results.¹¹ The sharp decrease in the electron number density after transition from laminar to turbulent flow may also be noticed.

In Fig. 2, the wake half-thickness for the 200,000-ft-alt case is shown. The results for the 150,000-ft alt are qualitatively similar. The inclusion of pressure gradient decreases the wake thickness. It is noticed that the initial thickness for the pressure gradient case is less than that for constant pressure. This is a consequence of applying the momentum theorem to relate the drag of the missile to the momentum and pressure defects in the wake. It then follows (see Ref. 1 for details) from the resulting expression that an increase (over ambient) in the pressure at the starting point results in a decrease in the starting thickness.

Although numerical results are not presented here, the calculations have shown that the centerline velocity and temperature are not significantly affected by pressure gradient.

In conclusion, the results have shown that pressure gradient has a significant effect on the wake thickness and electron number density, but only a small effect on fluid velocity and temperature. Specifically, pressure gradient decreases the wake thickness and reduces the electron number density. When these two effects are combined, it may be stated that a pressure gradient produces a much smaller volume (and number) of ionized particles than does the constant pressure case.

References

- ¹ Ness, N. and Fanucci, J. B., "Pressure gradient effects on nonequilibrium laminar and turbulent far wakes," Radio Corporation of America, Missile and Surface Radar Div., Down-Range Anti-Missile Measurement Program TM 63-13 (December 1963).
- ² Lees, L. and Hromas, L., "Turbulent diffusion in the wake of a blunt-nosed body at hypersonic speeds," J. Aerospace Sci. 29, 976-993 (1962).
- ³ Bloom, M. H. and Steiger, M. H., "Hypersonic axisymmetric wakes including effects of rate chemistry," General Applied Science Labs., TR-180 (September 1960).
- ⁴ Pallone, A. J., Erdos, J. I., and Eckerman, J., "Hypersonic laminar wakes and transition studies," Avco/RAD TM 63-33 (November 1963).
- ⁵ Steiger, M. H. and Bloom, M. H., "A note on hypersonic axisymmetric laminar wakes including rate chemistry and streamwise pressure gradients," General Applied Science Labs., TR-302 (August 1962).
- ⁶ Bortner, M. H., "Chemical kinetics in a reentry flow field," General Electric, TIS-R63SD63 (August 1963).
- ⁷ Ting, L. and Libby, P., "Fluid mechanics of axisymmetric wakes," General Applied Science Labs., TR-145A (June 1960); also "Remarks on the eddy viscosity in compressible moving flows," J. Aerospace Sci. 27, 797-798 (1960).
- ⁸ Zeiberg, S. L., "Correlation of hypersonic wake transition data," General Applied Science Labs., TR-382 (October 1963).
- ⁹ Sakurai, A., "On the propagation and structure of the blast wave, II," J. Phys. Soc. Japan 9, 256-266 (March-April 1954).
- ¹⁰ Lykoudis, P. S., "Laminar hypersonic trail in the expansion-conduction region," AIAA J. 1, 772-775 (1963).
- ¹¹ Lees, L., "Hypersonic wakes and trails," AIAA J. 2, 417-428 (1964); see Fig. 9.

Technical Comments

Comment on "Derivation of Element Stiffness Matrices"

G. C. BEST*

General Dynamics, Fort Worth, Texas

In a recent note, Pian¹ showed that the stiffness matrix of an element can be represented by

$$k = K_{aa} - K_{ab}K_{bb}^{-1}K_{ba} \quad (1)$$

where k is an $n \times n$ matrix, the development being based upon an assumption involving $n + l$ constants, the excess l constants being adjusted to minimize the potential energy. It may be of interest that Pian's result can also be represented

(in his notation) for a supported element by the simple formula

$$k = (B G^{-1} B^T)^{-1} \quad (2)$$

This follows from the fact that if a nonsingular matrix K and its inverse F partition into

$$\begin{bmatrix} K_{aa} & K_{ab} \\ K_{ba} & K_{bb} \end{bmatrix} \quad \text{and} \quad \begin{bmatrix} F_{aa} & F_{ab} \\ F_{ba} & F_{bb} \end{bmatrix} \quad (3)$$

respectively, then

$$F_{aa}^{-1} = K_{aa} - K_{ab}K_{bb}^{-1}K_{ba} \quad (4)$$

this being a well-known relation much used for eliminating coordinates from stiffness matrices.

It should be noted that Pian's matrix K inverts into $M^{-1}G^{-1}M^{-1T}$ and that

$$M^{-1} = \begin{bmatrix} B_a & B_b \\ 0 & I \end{bmatrix} \quad (5)$$

hence $F_{aa} = B G^{-1} B^T$, which gives Eq. (2) on inverting.

Received April 14, 1964.

* Senior Structures Engineer.

Furthermore, Eq. (2) can be derived in a very elementary way by adding extra rows B_c to B and extra coordinates q_c to the deflection vector q , so that Pian's Eq. (6) becomes

$$q^* = \begin{bmatrix} q \\ q_c \end{bmatrix} = \begin{bmatrix} B \\ B_c \end{bmatrix} \alpha = B^* \alpha \quad (6)$$

This can always be done so that B^* is nonsingular.

Then, by inverting Eq. (10) of Ref. 2, the flexibility matrix can be written

$$F = B^* G^{-1} B^{*T} \quad (7)$$

with F satisfying the equation

$$F Q^* = q^* \quad (8)$$

Now consider that certain forces Q^* are applied to the element causing deflections q^* , and then the forces at the newly added coordinates are removed, but the deflections q maintained. The deflections q_c , at the released coordinates, will adjust so as to minimize the strain energy, and by Eqs. (3) and (4), Eq. (7) reduces, on inversion, to Eq. (2).

References

¹ Pian, T. H. H., "Derivation of element stiffness matrices," AIAA J. 2, 576-577 (1964).

² Gallagher, R. H., "Techniques for the derivation of element matrices," AIAA J. 1, 1431-1432 (1963).

N-Segment Least-Squares Approximation

R. ARIS* AND M. M. DENN†

University of Minnesota, Minneapolis, Minn.

THE question of the best piecewise approximation to a given function by linear segments¹⁻³, and by step functions⁴ has been considered recently. It may be worthwhile to call attention to a property of the least-square criterion of fit which allows a rather easy calculation of the break points. Let $f(x)$, $a \leq x \leq b$, be approximated by N segments of a function of fixed form $\phi(x; c)$, where $c = (c_1, \dots, c_k)$ is a vector of constants to be chosen in each segment. The approximation will be best in the least-square sense if the Nk constants c_n , $n = 1, \dots, N$, and the $(N-1)$ break points x_n , $n = 1, \dots, N-1$, where $x_0 = a < x_1 < \dots < x_{N-1} < x_N = b$, are so chosen as to minimize

$$F_N(a, b) = \sum_{n=1}^N \int_{x_{n-1}}^{x_n} [f(x) - \phi(x; c_n)]^2 dx \quad (1)$$

This criterion of fit has the property that, over intervals where $f(x)$ is continuous, the segments are either continuous or, at a point of discontinuity, are equidistant from the curve $f(x)$. For if the break points x_{n-1} , x_n are chosen, the best value of the constants c_n satisfy the equations

$$\int_{x_{n-1}}^{x_n} [f(x) - \phi(x; c)] (\nabla_c \phi(x; c)) dx = 0 \quad (2)$$

where $\nabla_c \phi$ denotes the vector of partial derivatives of ϕ with respect to the components of c . Since x_n occurs once as an upper limit and once as a lower limit, the derivative of the

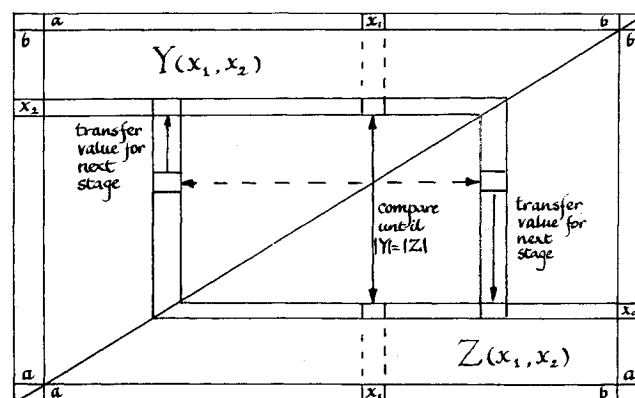


Fig. 1 Tabular method of generating best break points.

right-hand side of (1) with respect to x_n becomes, after using (2),

$$[f(x_n) - \phi(x_n; c_n)]^2 - [f(x_n) - \phi(x_n; c_{n+1})]^2$$

If x_n is chosen to minimize F_N , we have either

$$\phi(x_n; c_n) = \phi(x_n; c_{n+1})$$

or

$$f(x_n) = \frac{1}{2} \{ \phi(x_n; c_n) + \phi(x_n; c_{n+1}) \}$$

The value of this property may be illustrated by a tabular procedure for $N = 2$; it will be evident that a graph could be used in a similar way. Let the best fitting single segment for the interval (y, z) , $a < y < z < b$, be calculated. The constants $c(y, z)$ and the fit $F_1(y, z)$ may be recorded in double-entry table. Let

$$Y(y, z) = f(y) - \phi[y; c(y, z)]$$

$$Z(y, z) = f(z) - \phi[z; c(y, z)]$$

be tabulated as in Fig. 1. In the lower right-hand triangle, the deviation at the upper end of a single segment fit to the interval (x_0, x_1) is entered in a row corresponding to x_0 and column x_1 . In the upper left-hand triangle, the entry $Y(y_1, x_2)$ is in a row corresponding to x_2 and column x_1 . To find the best break point for two segments, we have only to compare entry of Y in the x_2 row with that of Z in the x_0 row and the same column in order to see at what value of x_1 we have either $Y = Z$ or $Y = -Z$. If more than one possibility exists, the best choice may easily be made by reference to the table of F_1 . With a sufficient number of entries, linear interpolation should suffice.

The process can be continued indefinitely by replacing the lower right-hand triangle by a table of $Z_2(x_0, x_2)$, the deviation at the upper end of the best two-segment fit. But this is merely to move the entry $Z(x_1, x_2)$ to the row x_0 as indicated by the arrows. If the same thing is done in the upper right-hand triangle to give $Y_2(x_2, x_4)$, the lower-end deviation of the best two-segment fit over the interval (x_2, x_4) , the table may now be used for calculating the best four-segment fit. If $N = 2^p$, only p repetitions of the process are required. Of course, record would be kept of the break points at each stage, and then the best constants could be obtained immediately from the single-segment tables.

References

¹ Aris, R., *Discrete Dynamic Programming* (Blaisdell, New York, 1964).

² Bellman, R. E. and Dreyfus, S. E., *Applied Dynamic Programming* (Princeton University Press, Princeton, N. J., 1962).

³ Bellman, R. E. and Kotkin, B., "On the approximation of curves by line segments using dynamic programming—II," Rand Corp. RM-2978-PR (1962).

⁴ Lubowe, A. G., "Optimal functional approximation using dynamic programming," AIAA J. 2, 376-377 (1964).

⁵ Stone, H., "Approximation of curves by line segments," Math. Computation 15, 40-47 (1961).

Received May 1, 1964.

* Professor, Department of Chemical Engineering.

† Graduate Student, Department of Chemical Engineering.